

Pulling a Mass

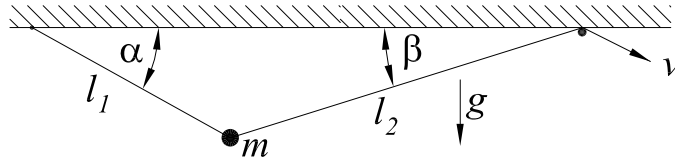
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1 Problem

Mass m is suspended from a fixed point A by a string of fixed length l_1 that makes angle α to the horizontal, while also pulled at constant speed v by a second string that passes over a fixed pulley. What is the tension T in the second string when that string makes angle β to the horizontal, as in the sketch below?



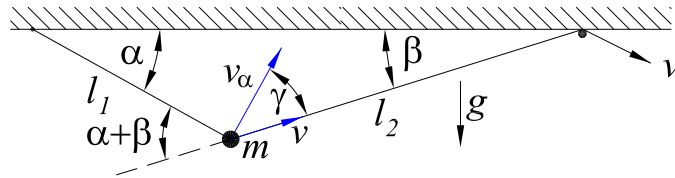
2 Solution

This problem is something of a “trick” question as a simpler answer can be given in terms of two angles α and β , although there is only one independent angle, say, α . It was posed as Prob. 21 on p. 188 of [1], with a solution by Newtonian methods given on p. 208. Here we give a solution via Lagrange’s method, which requires use of a Lagrange multiplier as there exists a velocity-dependent (nonholonomous) constraint.

The velocity constraint can usefully be expressed in terms of angles α and β by noting that the azimuthal velocity $v_\alpha = -l_1\dot{\alpha} = -l_1 d\alpha/dt$ of mass m is related to the velocity v of mass m along the direction of string 2 by

$$v = v_\alpha \cos \gamma = v_\alpha \sin(\alpha + \beta) = -l_1 \dot{\alpha} \sin(\alpha + \beta), \quad (1)$$

since angle $\gamma = \pi/2 - (\alpha + \beta)$, referring to the figure below.



The prescription¹ to include the effect of the constraint (1) in Lagrange’s equation is that

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\alpha}} - \frac{\partial \mathcal{L}}{\partial \alpha} = -l_1 \sin(\alpha + \beta) \lambda, \quad (2)$$

where λ is the Lagrange multiplier, with the physical significance of the tension T in string 2. Note that $Tl_1 \sin(\alpha + \beta)$ is the torque due to tension T about the support point of string 1, so the right side of eq. (2) is the generalized force corresponding to angle α .

¹Notes by the author on the method of Lagrange multiplier are around p. 68 of [2].

In the present problem, the Lagrangian \mathcal{L} based on independent angle α is

$$\mathcal{L} = \text{KE} - \text{PE} = \frac{ml_1^2 \dot{\alpha}^2}{2} + mgl_1 \sin \alpha, \quad (3)$$

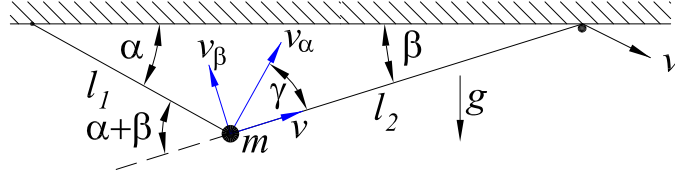
so eq. (2) implies that

$$ml_1^2 \ddot{\alpha} - mgl_1 \cos \alpha = -l_1 \sin(\alpha + \beta)T, \quad T = \frac{mg \cos \alpha}{\sin(\alpha + \beta)} - \frac{ml_1 \ddot{\alpha}}{\sin(\alpha + \beta)}. \quad (4)$$

A relation for the acceleration $a_\alpha = -l_1 \ddot{\alpha}$ of mass m with respect to the fixed point of string 1 can be found by first noting that the radial acceleration (with respect to the fixed point of string 1) has magnitude

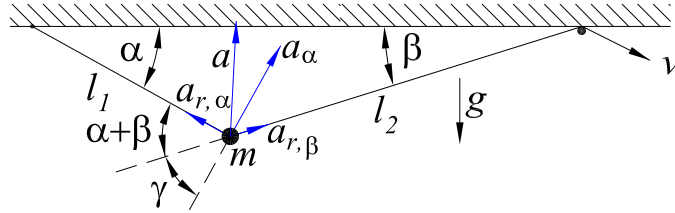
$$a_{r,\alpha} = \frac{v_\alpha^2}{r} = \frac{v^2}{l_1 \sin^2(\alpha + \beta)}. \quad (5)$$

A subtle trick is to consider also the radial acceleration $a_{r,\beta} = v_\beta^2/l_2$ of mass m with respect to the fixed pulley, as sketched in the two figures below.



We have that $v_\beta = v_\alpha \sin \gamma = v_\alpha \cos(\alpha + \beta)$, and hence the magnitude of the radial acceleration with respect to the pulley is

$$a_{r,\beta} = \frac{v_\beta^2}{l_2} = \frac{v_\alpha^2 \cos^2(\alpha + \beta)}{l_2} = a_\alpha \cos \gamma - a_{r,\alpha} \cos(\alpha + \beta) = a_\alpha \sin(\alpha + \beta) - a_{r,\alpha} \cos(\alpha + \beta). \quad (6)$$



From eq. (6) we have, recalling eqs. (1) and (5) and noting that $l_1/\sin \beta = l_2/\sin \alpha$,

$$\begin{aligned} a_\alpha = -l_1 \ddot{\alpha} &= \frac{a_{r,\beta}}{\sin(\alpha + \beta)} + \frac{a_{r,\alpha} \cos(\alpha + \beta)}{\sin(\alpha + \beta)} = \frac{v_\alpha^2 \cos^2(\alpha + \beta)}{l_2 \sin(\alpha + \beta)} + \frac{v^2 \cos(\alpha + \beta)}{l_1 \sin^2(\alpha + \beta)} \\ &= \frac{v^2 \cos^2(\alpha + \beta)}{l_2 \sin^3(\alpha + \beta)} + \frac{v^2 \cos(\alpha + \beta)}{l_1 \sin^2(\alpha + \beta)}. \end{aligned} \quad (7)$$

Finally, from eq. (4), the tension in string 2 is²

$$T = \frac{mg \cos \alpha}{\sin(\alpha + \beta)} + \frac{mv^2 \cot(\alpha + \beta)}{\sin^2(\alpha + \beta)} \left(\frac{\cot(\alpha + \beta)}{l_2} + \frac{1}{l_1} \right). \quad (8)$$

²There is a small error in the expression for T on p. 209 of [1].

We can use eq. (8) for the tension T in the first of eq. (4) to obtain an equation of motion involving both angles α and β . Since angle β is a known function of angle α , we could display a (complicated) equation of motion in terms of angle α only.

The author was introduced to this problem by Clement Chan.

References

- [1] J. Wang and B. Ricardo, *Competitive Physics: Mechanics and Waves*, (World Scientific, 2019). http://kirkmcd.princeton.edu/examples/mechanics/wang_19.pdf
- [2] K.T. McDonald, *Physics 205 Lecture 6*.
<http://kirkmcd.princeton.edu/examples/Ph205/ph205l6.pdf>